

## Chemical Evolution of the Universe

### Problem sheet 2

1. (a) In the lecture we derived a relationship between the Hubble parameter  $H$  and the temperature  $T$  for the very early Universe where the mass-energy density is dominated by relativistic particles. Using your knowledge of  $R(t)$  during this period from problem sheet 1, show that

$$kT(t) = \frac{0.85 \text{ MeV}}{\sqrt{t/s}}$$

- (b) How does this relation change if photons are the only relativistic particles (instead of photons,  $e^-$ ,  $e^+$ ,  $\nu_e$  and  $\bar{\nu}_e$ , as assumed in the lecture)?

**3 points**

2. As we have seen in the lecture, nucleosynthesis of  ${}^4\text{He}$  and heavier nuclei requires the existence of deuterium first (deuterium bottleneck).

- (a) Compute the binding energy of a deuterium nucleus.  
 (b) What is the frequency of a photon capable of destroying a deuterium nucleus?  
 (c) At what temperature does a black body emit most of its photons at the above frequency?  
 (d) At what time does the Universe reach this temperature (using the relation from 1.(b))?

**4 points**

3. Contrary to the result above, your lecture notes will tell you that deuterium formation sets in somewhat later. We will now learn why.

- (a) The present-day density parameter of baryons is  $\Omega_{b,0} = 0.045$ . Let us assume that all baryons are protons and that  $H_0 = 70 \text{ km/s/Mpc}$ . Calculate the present-day number density of baryons,  $n_b$ .  
 (b) The present-day temperature of the Cosmic Microwave Background is  $T_0 = 2.73 \text{ K}$ . Use the Bose-Einstein distribution to compute the present-day number density of photons,  $n_\gamma$ .  
 Hint:  $\int_0^\infty \frac{x^2}{e^x - 1} dx = 2.4041$   
 (c) What is the present-day value of the baryon-to-photon ratio,  $\eta = \frac{n_b}{n_\gamma}$ ?  
 (d) How does  $\eta$  evolve with time (qualitatively)? Specifically, at the time you found in 2.(d), is  $\eta$  larger or smaller than its present-day value, or equal to it? Give your reasoning.

Let's summarise what we've found out so far. At the time in 2.(d), a typical photon has just about enough energy to dissociate a deuterium nucleus. At later times, when the Universe has further cooled, a typical photon will therefore *not* have enough energy to dissociate a deuterium nucleus. However, in 3.(d) we found out that there are *many* more photons than baryons in the Universe at this time. Thus, although a *typical* photon cannot dissociate a deuterium nucleus any longer, there are still plenty of higher energy photons that can.

- (e) At the time you found in 2.(d), compute the fraction of photons with energy larger than the binding energy of deuterium, i.e.  $n_\gamma(h\nu > \text{BE}_D)/n_\gamma$ , again using the Bose-Einstein distribution.  
 Hint 1: Use Wien's approximation (why is it ok to do that?)  
 Hint 2:  $\int x^2 e^{-ax} dx = \frac{e^{-ax}}{a^3} [-ax(ax+2) - 2]$   
 (f) Now calculate  $n_\gamma(h\nu > \text{BE}_D)/n_b$ .  
 (g) Repeat the calculation for  $t = 270 \text{ s}$  and verify that deuterium will start forming around this time.

**8 points**